

Connect ~~with~~

Hubble expansion
with Galaxy Counts

We have a conservation law

$$\int R^3 = \text{const}$$

If we \div by the mass of the galaxy this reads $nR^3 = \text{const}$

or $n = \alpha R^{-3}$ [α is the const]

$$\text{then } \frac{dn}{dt} = -3\alpha R^{-4} \frac{dR}{dt}$$

$$\left| \frac{\dot{n}}{n} \right| = + \frac{3\alpha R^{-4} \dot{R}}{\alpha R^{-3}}$$

$$|\dot{n}/n| = + 3 \dot{R}/R = + 3 H$$

[H is the Hubble const]

Let's write H as γt .

[the Hubble $\xrightarrow{\text{approx.}} H$.]

Then, from previous work in Problem set 1, $|\dot{n}/n| = 3.4 \text{ billion years} = t_H/3$
 $t_H = 3 \times 3.4 = 10.2 \text{ billion years}$

"Hubble's law cubed" is NOT universal

Suppose we observe 6 galaxies: Galaxy A, B, C, D, E and F. We measure their distances D (in megaparsecs) and velocities from us V (in kilometers / second). Here's what we find

HUBBLE'S COSMIC TABLE

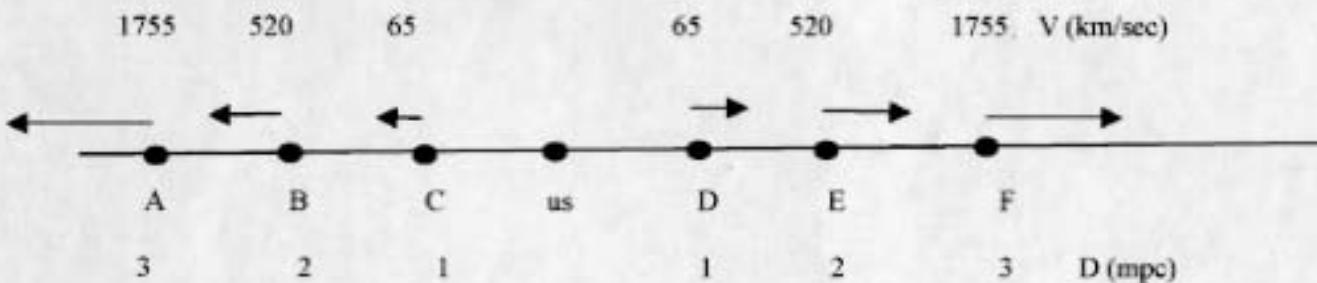
Galaxy	D	V
A	3	1755
B	2	520
C	1	65
D	1	65
E	2	520
F	3	1755

Now we do a little math: we divide V by D³ and we get 65 in every case. So we conclude

$$V = H D^3$$

Where H = 65. This, of course, is not Hubble's law.

Here's a diagram illustrating what we have found:



Now we want to ask a new question: what would an alien cosmologist living on another galaxy find? To be specific, let's suppose that our alien's name is Elbbuh, and it (they don't have sexes) lives in galaxy D. We need to find

- (a) the *distances Elbbuh measures*. These are the distances of each galaxy from D. They are

Galaxy	Distance from galaxy D (mpc)
A	4
B	3
C	2
us	1
E	1
F	2

(b) the *velocities Elbbuh measures*. These are the velocities of all the galaxies relative to D.

These are a little harder to find. Let's start with "us" — this is pretty easy. Since D is going 65 km/sec away from us, we are going 65 km/sec away from it.

Now let's study galaxy C. In our above diagram, it is going 65 km/sec to the left, and D is going 65 km/sec to the right. So the velocity of C relative to D is $65+65=130$ km/sec.

Similarly, in the above diagram, E is going 520 km/sec to the right, and D is going 65 km/sec to the right. D is "trying to keep up with E but not going fast enough to do so." So the velocity of E relative to D is $520 - 65 = 455$ km/sec.

We keep on in the same way, and finally reach the following results

ELBBUH'S COSMIC TABLE

Galaxy	Distance from galaxy D (mpc)	Velocity relative to galaxy D (km/sec)	Velocity divided by D cubed
A	4	1820	28.4
B	3	585	21.7
C	2	130	16.3
us	1	65	65
E	1	455	455
F	2	1690	211

Now here's the final step. Let's divide the velocity by the distance cubed to find "Elbbuh's law." The results are shown in the right-hand column of the above table. Notice that one of the entries is 65 — but the others are not. This proves that

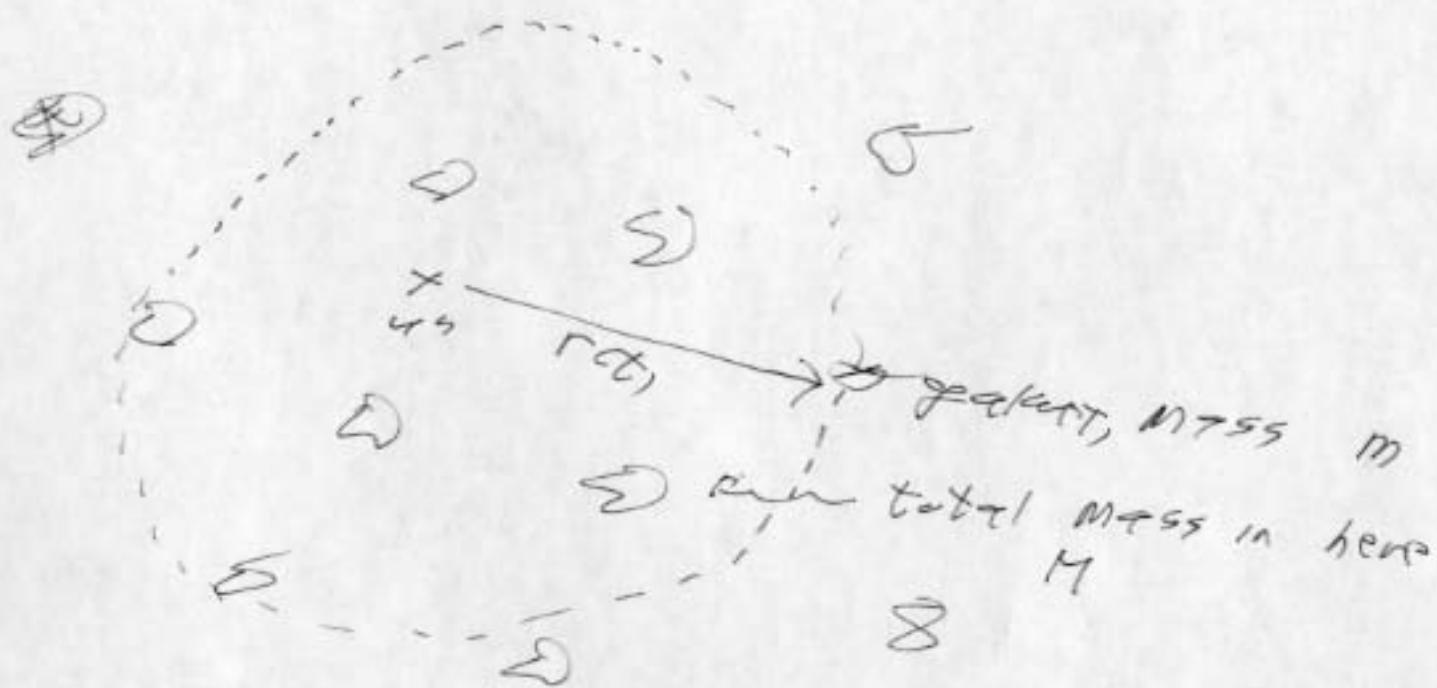
Velocity_{relative to D} does NOT equal $(65)(\text{Distance}_{\text{relative to D}})^3$

So Elbbuh's law is not the same as Hubble's law. This proves that '*Hubble's law cubed*' is not universal.

(3) Friedman ~~exp~~ eqs in the early universe

(4)

(*) we proceed just like in classifying the force of gravity on the galactic.



$$\text{The force is } F = \vec{r} \frac{GMm}{r^2}$$

because the force decelerates the galaxy's motion

$$M = \frac{4}{3}\pi r^3 \rho \quad \text{so } F = -\frac{4}{3}\pi Gm\rho r$$

$$\text{and } F = ma = m \ddot{r} \quad \text{so}$$

$$\ddot{r} = -\frac{4}{3}\pi a \rho r$$

in terms of the
expansion parameter:

(5)

$$\Gamma(t) = R(t) \Gamma(t_0)$$

$$\ddot{\Gamma}(t) = \ddot{R}(t) \Gamma(t_0)$$

so we get

$$\ddot{R}(t) = -\frac{4}{3}\pi a \rho(t) R(t)$$

This is exactly what we did in class.
Here's the new twist: write

$$\rho(t) = \frac{\text{const}}{R^4(t)} \quad \text{where the}$$

constant is $\rho(t_0) R^{12}(t_0) = \rho_0$

$$\underbrace{\qquad}_{=1}$$

so we get

$$\ddot{R}(t) = -\frac{4}{3}\pi a \rho(t_0) / R^3(t)$$

new
1st
eq

Multiply by \dot{R}

$$\ddot{R} \dot{R} = -\frac{4}{3}\pi a \rho(t_0) \frac{\dot{R}}{R^3}$$

\curvearrowleft

A

B

A is the derivative $\Rightarrow \frac{1}{2} \dot{R}^2$ (6)

B is the derivative of $\left[-\frac{1}{2} R^{-2}\right]$

so the above equation can be written

$$\frac{d}{dt} \left[\frac{1}{2} \dot{R}^2 \right] = \frac{d}{dt} \left[-\frac{4}{3} \pi \alpha \rho(t) \right] \left[-\frac{1}{2} R^{-2} \right]$$

so

$$\frac{1}{2} \dot{R}^2 = \left[-\frac{4}{3} \pi \alpha \rho(t_0) \right] \left[-\frac{1}{2} R^{-2} \right] + \text{const}$$

$$\dot{R}^2 = \frac{4}{3} \pi \alpha \rho(t_0) \frac{1}{R^2} + \text{const}$$

2nd new ~~Friedmann~~ eq

(b) Let's verify that $R(t) = \alpha t^{1/2}$ solves the equation if "const" is zero.

$$\frac{dR}{dt} = \frac{d}{dt} [\alpha t^{1/2}] = \frac{1}{2} \alpha t^{-1/2}$$

is the left-hand side

$$\text{so } \dot{R}^2 = \frac{1}{4} \alpha^2 t^{-1}$$

The right hand side is

$$\frac{\frac{4}{3}\pi\alpha^3(t)}{[\alpha t^{\frac{1}{2}}]^2} = \frac{4}{3}\pi\alpha^3(t) / t^2$$

equating

$$\frac{1}{4}\alpha^2 t^{-1} = \frac{4}{3}\pi\alpha^3(t) / t^2$$

the t 's cancel, and if

$$\alpha^4 = \frac{16}{3}\pi\alpha^3(t) \quad \text{the}$$

two sides are equal.

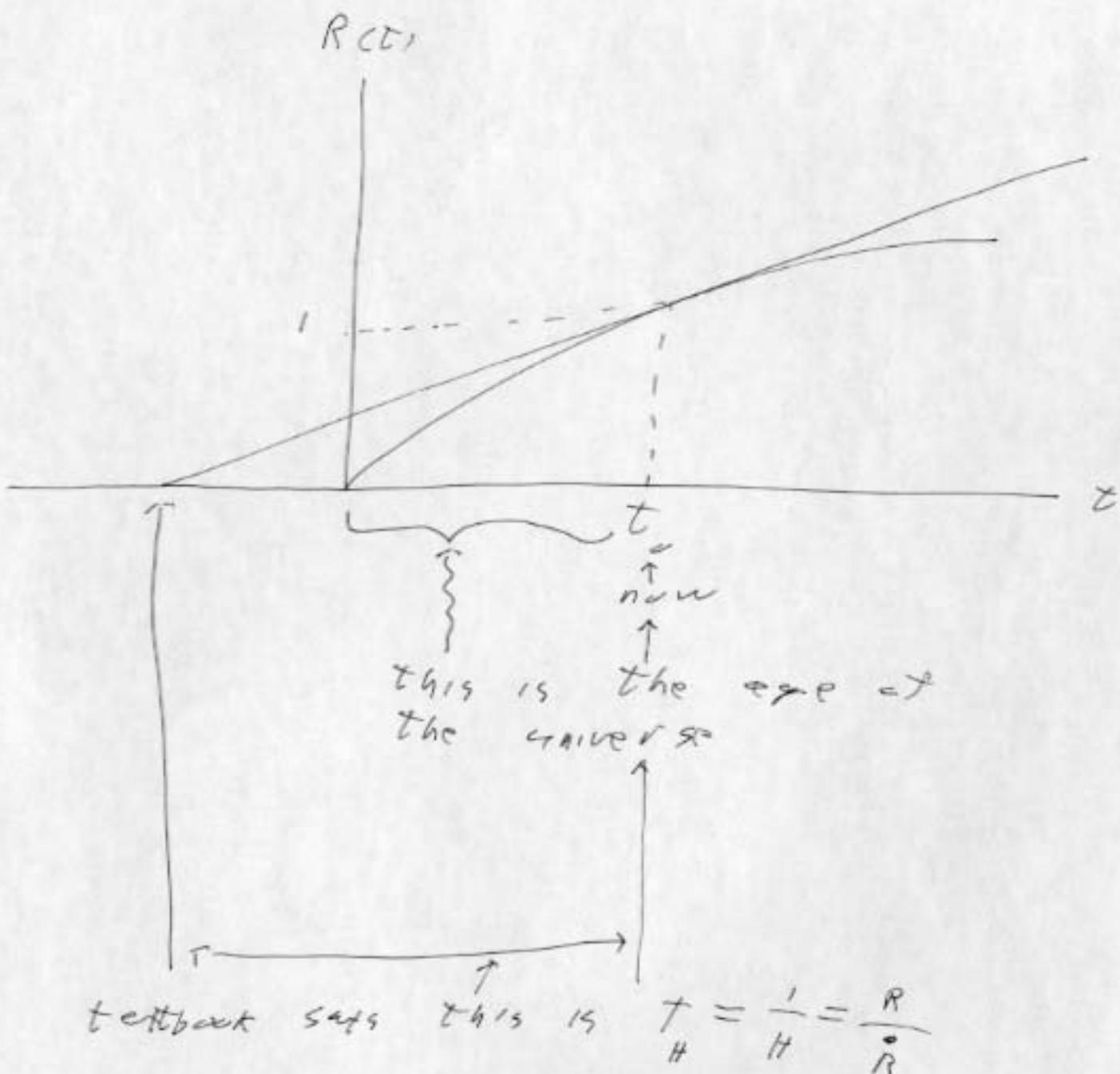
$$\alpha = \left[\frac{16}{3}\pi\alpha^3(t) \right]^{\frac{1}{4}}$$

(c) yes, there was a big bang, since
~~at~~ when $t=0$ R was zero and
the density was infinite

4 The Hubble Law

(8)

I can think of 2 ways to do
this. Either is fine.



but why?

one way: intuitively

The ~~redshift~~ has constant slope, and

$$\text{dist to } z \text{ is } \propto \text{dist to } 1 \text{ au}$$
$$r(t) = r(t_0) R(t)$$

$$\text{so } \frac{dr}{dt} = v(t) = v(t_0) R$$

If R is constant, $v(t)$ is too

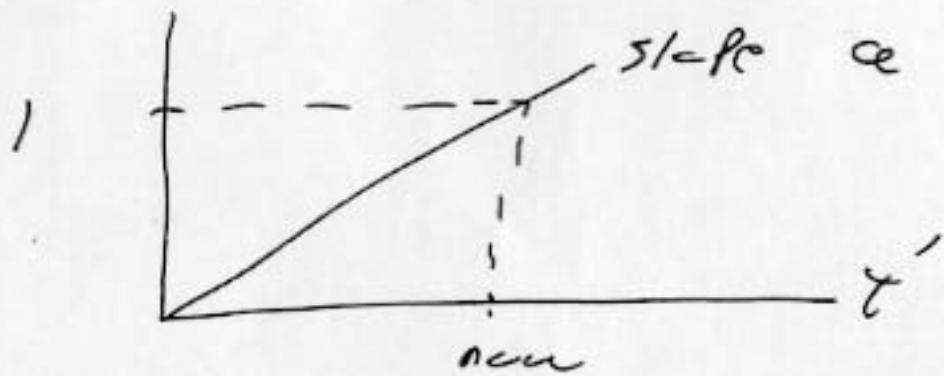
and we proved in class that, if galaxies move with constant velocity, the size \rightarrow the inverse

$$\text{is } T_H = \frac{1}{H}$$

second way: mathematically

let's first draw a graph \rightarrow the straight line ~~is~~ \propto a different time axis t'

$f(t')$



(10)

(10)

The eq for the line is $f(t') = \alpha t'$
and it has the same 1/slope

$$t' = \frac{1}{\alpha} = \frac{1}{\text{slope}}$$

So, on the graph on the previous page,
the interval of time ~~is~~ which
the textbook calls τ_H is $\frac{1}{\text{slope}}$
 $= \frac{1}{R}$

But we can multiply this

by $R(t_0)$ since this = 1, so we get
 $\tau_H = \frac{R}{R} \Big|_{t=t_0}$ but this is the res.
it is the definition of τ_H

(5) At big value of H

(11)

(11) the distance to the sun (Appendix) is

$$1 \text{ AU} = 1.50 \times 10^{13} \text{ cm}$$

so the circumference of our orbit is $2\pi \times 1.50$.
The earth covers this distance in 1 year, so
our velocity is

$$v = \frac{2\pi(1.5 \times 10^{13}) \text{ cm}}{365 \times 24 \times 60 \times 60 \text{ sec}} = 2.77 \times 10^6 \frac{\text{cm}}{\text{sec}}$$

and we are proposing that the sun is
receding at $\frac{1}{100}$ this speed, or $2.77 \times 10^4 \frac{\text{cm}}{\text{sec}}$.
Setting this = $H \times 1.4$. we solve for
 H .

$$H = \frac{2.77 \times 10^4 \text{ cm/sec}}{1.4} = 1.99 \times 10^{-7} \cancel{\text{sun}}^{-1}$$

$$\frac{1}{H} = 15.9 \text{ years !}$$

