The homotopy moment map for Chern–Simons field theory

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Homotopy moment maps

Let *X* be a manifold, with de Rham complex $\Omega^*(X)$. Let *n* be a natural number.

Definition

An *n*-multisymplectic form on X is a closed differential form $\omega \in \Omega^{n+1}(X)$ such that the morphism

$$\xi \in \operatorname{Vec}(X) \mapsto \iota(\xi)\omega \in \Omega^n(X)$$

is injective.

The case n = 1 of symplectic forms is special, since if X is finite-dimensional, the morphism $\xi \mapsto \iota(\xi)\omega$ is not just injective, but bijective.

Let \mathfrak{h} be a Lie algebra with differential action on X:

$$\rho:\mathfrak{h}\to \operatorname{Vec}(X).$$

Let $C^*(\mathfrak{h})$ be the Chevalley–Eilenberg complex of \mathfrak{h} , with differential δ_0 . Denote the tensor product of $C^*(\mathfrak{h})$ and $\Omega^*(X)$ by

 $C^*(\mathfrak{h})\otimes \Omega^*(X).$

Choose a basis $\{c_i\}$ of \mathfrak{h} , and dual basis $\{c^i\}$ of \mathfrak{h}^* ; denote the structure coefficients of \mathfrak{h} by A_{ii}^k . Form the operator

$$\iota = \sum_{i} c^{i} \otimes \iota(\rho_{i}) : C^{p}(\mathfrak{h}) \otimes \Omega^{q}(X) \to C^{p+1}(\mathfrak{h}) \otimes \Omega^{q-1}(X),$$

of total degree 0.

Definition (Callies, Frégier, Rogers, Zambon)

A **homotopy moment map** for the action ρ of \mathfrak{h} on an *n*-multisymplectic manifold (X, ω) is an element $\mu \in C^*(\mathfrak{h}) \otimes \Omega^*(X)$ of total degree n - 1 such that

$$(\delta_0 + d)\mu = e^t\omega - \omega.$$

The case of symplectic manifolds: n = 1

If n = 1, we have $\iota \omega = \sum_i c^i \iota(\rho_i) \omega \in \mathfrak{h}^* \otimes \Omega^1(X)$. Thus, a homotopy moment map is an element $\mu = \sum_i c^i \mu_i \in \mathfrak{h}^* \otimes \Omega^0(X)$ such that

$$d\mu_i = -\iota(\rho_i)\omega$$

and

$$\sum_{k} \mathbf{A}_{ij}^{k} \mu_{k} = \iota(\rho_{i})\iota(\rho_{j})\omega = -\mathcal{L}(\rho_{i})\mu_{j} = \{\mu_{i}, \mu_{j}\}.$$

In other words, μ is a moment map for the action ρ of \mathfrak{h} on the symplectic manifold (X, ω) .

Reformulation for exact *n*-multisymplectic manifolds

Definition

An *n*-multisymplectic manifold (X, ω) is **exact** if $\omega = d\alpha$, for $\alpha \in \Omega^n(X)$.

The Lie derivative $c_i \mapsto \mathcal{L}(\rho_i)$ makes the de Rham complex $\Omega^*(X)$ into a differential graded \mathfrak{h} -module. This motivates replacing the complex $C^*(\mathfrak{h}) \otimes \Omega^*(X)$ by the complex $C^*(\mathfrak{h}, \Omega^*(X))$, with differential $\delta + d$, where

$$\delta = \delta_0 + \sum_i c^i \mathcal{L}(\rho_i).$$

Lemma

$$\delta + d = e^{\iota} \circ (\delta_0 + d) \circ e^{-\iota}$$

The equation for a homotopy moment map becomes

$$(\delta_0 + d)(\alpha + \mu) = e^{\iota}\omega.$$

Define

$$v = e^{\iota}(\alpha + \mu) \in C^*(\mathfrak{h}, \Omega^*(X)),$$

or equivalently,

$$\mu = \mathbf{e}^{-\iota} \mathbf{v} - \alpha.$$

For exact *n*-multisymplectic manifolds, a homotopy moment map is an equivariant extension of α :

$$(\delta + d)v = \omega.$$

The variational bicomplex

Let $p : E \to M$ be a fiber bundle, with jet-space $J_{\infty}(p)$. An **adapted** coordinate system at a point $e \in E$ is a coordinate system $(t^1, \ldots, t^n; u^1, \ldots, u^N)$ such that the coordinates (t^1, \ldots, t^n) are pulled back from a coordinate system around $p(e) \in M$.

The coordinates t^{μ} are the **independent coordinates**, and *M* is the **worldsheet**; the coordinates u^a are the **dependent coordinates**, identified with the fields of the theory. Denote $\partial/\partial t^{\mu}$ by ∂_{μ} .

An adapted coordinate system gives rise to coordinates

$$(t^1,\ldots,t^n;\partial^l u^1,\ldots,\partial^l u^N)$$

on the jet-space $J_{\infty}(p)$, where *I* ranges over multi-indices $(i_1, \ldots, i_n) \in \mathbb{N}^n$, and

$$\partial^{l} = \partial_{1}^{i_{1}} \dots \partial_{n}^{i_{n}}$$

Denote partial differentiation $\partial/\partial(\partial^l u_i)$ with respect to a jet coordinate by $\partial_{i,l}$.

Let O_{∞} be sheaf of algebras over M whose sections are smooth functions in the coordinates $(t^1, \ldots, t^n; \partial^l u^1, \ldots, \partial^l u^N)$.

The variation de Rham complex

The variational de Rham complex Ω_{∞}^* is the de Rham complex associated to O_{∞} .

It is bigraded: the differentials dt^{μ} have bidegree (1,0), and the differentials $\theta_{l}^{i} = d(\partial^{l}u^{i}) - \sum_{\mu} \partial_{\mu}\partial^{l}u^{i} dt^{\mu}$

have bidegree (0, 1).

The differential d on Ω_{∞} breaks into horizontal and vertical parts

$$d^{1,0} = \sum_{\mu} dt^{\mu} D_{\mu} : \Omega_{\infty}^{p,q} \to \Omega_{\infty}^{p+1,q}, \quad d^{0,1} = \sum_{i,l} \theta_l^a \, \partial_{i,l} : \Omega_{\infty}^{p,q} \to \Omega_{\infty}^{p,q+1}.$$

where

$$D_{\mu} = \partial_{\mu} + \sum_{i,l} \partial_{\mu} \left(\partial^{l} u^{i} \right) \partial_{i,l}.$$

The variational *n*-multisymplectic form of a classical field theory

A first-order Lagrangian density L determines a classical field theory. In an adapated coordinate system,

$$L = L(t, u_i, \partial_{\mu} u_i) dt^1 \wedge \ldots \wedge dt^n \in \Omega_{\infty}^{n,0}.$$

From *L*, we construct a Lepage form

$$\alpha = L - \partial_{i,\mu}L \,\theta_i \wedge \iota(\partial_{\mu}) \,dt^1 \wedge \ldots \wedge dt^n \in \Omega^{n,0}_{\infty} \oplus \Omega^{n-1,1}_{\infty} \subset \Omega^n_{\infty}.$$

The differential $\omega = d\alpha \in \Omega_{\infty}^{n+1}$ extends Noether's symplectic form off-shell:

$$\omega = \frac{\delta L}{\delta u_i} \, du_i \wedge dt^1 \wedge \ldots \wedge dt^n + \cdots \in \Omega_{\infty}^{n,1} + \Omega_{\infty}^{n-1,2}.$$

Modulo mild conditions on *L* which guarantee nondegeneracy, ω is a variational analogue of a *n*-multisymplectic form.

Variational Cartan calculus

Derivations $\xi \in \text{Der}(O_{\infty})$ of O_{∞} are the variational vector fields. There is a variational Cartan calculus: to a variational vector field ξ , we associate the contraction

$$\iota(\xi):\Omega^*_{\infty}\to\Omega^{*-1}_{\infty}$$

and the Lie derivative

$$\mathcal{L}(\xi) = \boldsymbol{d} \circ \iota(\xi) + \iota(\xi) \circ \boldsymbol{d} : \Omega_{\infty}^* \to \Omega_{\infty}^*.$$

Fix a fiber bundle $p : E \to B$ and a (first-order) Lagrangian density $L \in \Omega_{\infty}^{n,0}$ with associated Lepage form $\alpha \in \Omega_{\infty}^{n}$ and variational *n*-multisymplectic form $\omega \in \Omega_{\infty}^{n+1}$.

Consider a Lie algebra h and a variational action of h

$$\rho:\mathfrak{h}\to\mathsf{Der}(\mathcal{O}_\infty)$$

such that $\mathcal{L}(\rho)\omega = 0$.

Definition

A variational homotopy moment map for ρ is an element $v \in C^*(\mathfrak{h}, \Omega_{\infty}^*)$ of total degree n - 1 such that

 $(\delta + d)v = \omega.$

The Chern–Simons classical field theory

Let *M* be a closed oriented 3-manifold. Let g be a reductive Lie algebra, with invariant inner product (-, -). Let *G* be a connected compact Lie group with Lie algebra g.

Let P be a G-principal bundle over M. We consider a field theory over M, where the fields (dependent coordinates) are the components of a connection over P.

Over a chart $U \subset M$ where *P* is trivialized, the Lepage form of the Chern–Simons Lagrangian theory equals

$$\alpha = \frac{1}{2}(A, dA) + \frac{1}{6}(A, [A, A]).$$

If *G* is simply connected, then *P* is a trivial bundle, and the Lepage form is globally defined. In general, *P* is not trivial, and the Lepage form is a Čech cochain $\alpha \in \check{C}^*(\mathcal{U}, \Omega_{\infty}^*)$ of total degree 3.

The 3-multisymplectic form of Chern–Simons field theory

The 3-multisymplectic form equals

$$\omega = \frac{1}{2}(\mathbb{F}, \mathbb{F}) = (F, d^{0,1}A) + \frac{1}{2}(d^{0,1}A, d^{0,1}A),$$

where $F = d^{1,0}A + \frac{1}{2}[A, A]$ is the curvature, and

$$\mathbb{F} = dA + \frac{1}{2}[A, A] = d^{0,1}A + F$$

is its analogue in the variational bicomplex.

From this point, the slides from the talk contained an error, which is corrected here: the multisymplectic form of Chern–Simons theory is only nondegenerate on evolutionary vector fields.

This 3-form is nondegenerate on evolutionary vector fields: those of the form

$$\operatorname{pr}\left(X,\frac{\partial}{\partial A}\right) = \sum_{I} \left(\partial^{I} X, \frac{\partial}{\partial (\partial^{I} A)}\right),$$

where $X \in \Omega^{1,0}_{\infty} \otimes \mathfrak{g}$.

The Atiyah Lie algebra

The Atiyah Lie algebra $\mathcal{A}(P)$ of a principal bundle P is spanned by covariant derivatives ∇_{ξ}^{A} , and gauge transformations $\eta \in P \times_{G} \mathfrak{g}$.

In adapted coordinates associated to a local trivialization of *P* over an open subset $U \subset M$, we have $\nabla_{\xi}^{A} + \eta = \xi^{\mu}\partial_{\mu} + \operatorname{ad}(\iota(\xi)A + \eta)$.

The evolutionary vector fields

$$\rho\left(\nabla_{\xi}^{A} + \eta\right) = \Pr\left(-\mathcal{L}(\xi)A + \operatorname{ad} \nabla^{A}\left(\iota(\xi)A + \eta\right), \frac{\partial}{\partial A}\right)$$
$$= \Pr\left(-\iota(\xi)\mathbb{F} + \operatorname{ad} \nabla^{A}\eta, \frac{\partial}{\partial A}\right)$$

yield an action of the Atiyah Lie algebra on the variational de Rham complex.

The (local) homotopy moment map

We have
$$\iota \mathbb{F} = \nabla^A \eta$$
, and hence $e^{\iota} \mathbb{F} = \mathbb{F} + \nabla^A \eta$
Since $\delta_0 A = 0$, we see that $(\delta_0 + d)A + \frac{1}{2}[A, A] = \mathbb{F}$.
Applying e^{ι} to both sides, and using $e^{\iota}A = A$, we see that
 $(\delta + d)A + \frac{1}{2}[A, A] = e^{\iota}\mathbb{F}$.

Hence

$$(\delta+d)(A-\eta)+\tfrac{1}{2}[A-\eta,A-\eta]=\mathbb{F}.$$

Theorem

$$v = \frac{1}{2}(A - \eta, (\delta + d)(A - \eta)) + \frac{1}{6}(A - \eta, [A - \eta, A - \eta])$$

is a variational homotopy moment map for the Chern-Simons theory.

The (global) homotopy moment map

The point of this talk is that when P is not a trivial bundle, the Chern–Simons action, and Lepage form, are Čech cochains for a cover \mathcal{U} of M.

This forces the variational homotopy moment map

$$v \in \boldsymbol{C}^*(\boldsymbol{\mathfrak{g}}, \check{\boldsymbol{C}}^*(\boldsymbol{\mathcal{U}}, \Omega^*_\infty))$$

to be a Čech cochain of total degree 3, satisfying the equation

$$(\delta + \check{\delta} + d)v = \omega.$$

The formulas for the higher degree Čech cochains components of v are left as an exercise.

Since $\mu = e^{-\iota}v - \alpha$ is not an element of $C^*(\mathfrak{g}, \Omega^*_{\infty})$, we have an example of a homotopy moment map that is a differential character, but not a globally defined differential form.

This phenomenon only arises because n > 1.